



## EPR, Optical Absorption and Superposition Model Studies of Cr<sup>3+</sup> Doped Cesium Tetrabromozincate

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### Abstract

Electron paramagnetic resonance (EPR) study of Cr<sup>3+</sup> doped cesium tetrabromozincate (CTBZ) single crystal is done at room temperature. The hyperfine structure for Cr<sup>53</sup> isotope is also obtained. Two magnetically inequivalent sites for Cr<sup>3+</sup> are observed. The spin Hamiltonian parameters are evaluated as:  $D = 234 \times 10^{-4} \text{ cm}^{-1}$ ,  $E = 69 \times 10^{-4} \text{ cm}^{-1}$ ,  $g = 2.0104$ ,  $A = 80 \times 10^{-4} \text{ cm}^{-1}$  for site I and  $D = 235 \times 10^{-4} \text{ cm}^{-1}$ ,  $E = 70 \times 10^{-4} \text{ cm}^{-1}$ ,  $g = 2.0061$ ,  $A = 82 \times 10^{-4} \text{ cm}^{-1}$  for site II, respectively. The optical absorption spectra are recorded at room temperature. The energy values of different orbital levels are determined. The values of various parameters obtained are:  $B = 602 \text{ cm}^{-1}$ ,  $C = 2504 \text{ cm}^{-1}$ ,  $Dq = 1870 \text{ cm}^{-1}$ ,  $h = 1.63$  and  $k = 0.21$ , where B and C are Racah parameters, Dq is crystal field parameter, and h and k are nephelauxetic parameters, respectively. Theoretical zero-field splitting (ZFS) parameters for Cr<sup>3+</sup> at two sites in CTBZ are evaluated using superposition model and microscopic spin Hamiltonian theory. The theoretical ZFS parameters are in good agreement with the experimental values.

**Keywords:** A. Inorganic compounds; B. Crystal Growth; D. Crystal Fields; D. Electron Paramagnetic Resonance; D. Optical Properties.

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### 1. Introduction

Electron paramagnetic resonance (EPR) study of transition metal ion doped single crystals is used to find the local site symmetry and zero-field splitting parameters [1]. The optical study is used to obtain crystal field strength and the energy level structure of the impurity ion. Cr<sup>3+</sup> is one of the most investigated transition metal ions used as a probe for studying the structure and local symmetry of the crystal field in new materials [2–4]. Chromium is used as tanning agent in leather industry [5].

The theoretical studies on the spin Hamiltonian parameters of 3d ions in crystals suggested various mechanisms to contribute to ground state splitting of magnetic ion. The mostly used perturbation procedure treats cubic field and diagonal parts of free ion Hamiltonian as unperturbed Hamiltonian, whereas the spin-orbit coupling, the low-symmetry field, and off-diagonal part of free ion Hamiltonian are considered as perturbation [6].

Cesium tetrabromozincate, Cs<sub>2</sub>ZnBr<sub>4</sub> (CTBZ) can show thermo sensitive photo-selective phenomenon like ammonium tetrabromozincate [7]. This developed our interest in CTBZ. In the present investigation, EPR and optical studies of Cr<sup>3+</sup> ion in CTBZ single crystal are done to obtain information about the crystal field strength, energy level structure and the local site symmetry of the doped ion. The study is further used to describe the nature of bonding in the complex. Furthermore, zero-field splitting (ZFS) parameters are determined theoretically using superposition model (SPM) and microscopic spin Hamiltonian theory. The theoretical ZFS parameter values are in good agreement with the experimental ones.

## 2. Crystal Structure

The crystal structure of CTBZ is orthorhombic with space group  $Pnma$  having the unit cell dimensions  $a = 10.196 \text{ \AA}$ ,  $b = 7.770 \text{ \AA}$ ,  $c = 13.517 \text{ \AA}$ . The number of atoms per unit cell is four i.e.  $Z = 4$  [8]. The bromine atoms Br (1), Br (2), Br (3) and Br (4) are located at the distances of  $7.518 \text{ \AA}$ ,  $3.634 \text{ \AA}$ ,  $6.974 \text{ \AA}$  and  $6.974 \text{ \AA}$  from  $\text{Cs}^+$  ion in the host lattice.

## 3. Experimental

CTBZ single crystals were grown at room temperature by slow evaporation of an aqueous solution of cesium bromide and zinc bromide to which 0.1 wt% of  $\text{CrCl}_3 \cdot 6\text{H}_2\text{O}$  was added. The crystals grew in about 15-20 days. EPR spectra of  $\text{Cr}^{3+}$  doped CTBZ were recorded at room temperature on Varian X-Band E-112 (9.52 GHz) reflection type EPR spectrometer about the three mutually perpendicular crystallographic axes **a**, **b** and **c** at an interval of  $10^\circ$  each using a goniometer. The optical spectra of the crystal were recorded on Unicam-5625 spectrophotometer in UV-Vis-NIR (195–1100 nm) region at room temperature.

## 4. Results and Discussion

The recorded EPR spectrum of  $\text{Cr}^{3+}$  doped CTBZ single crystal at room temperature, when the external magnetic field is along the **c**-axis, is shown in Fig.1 (a). EPR spectra consist of 24 lines with hyperfine structure. These lines in all three planes show the presence of two magnetically inequivalent sites of  $\text{Cr}^{3+}$  ions. The hyperfine structure shows the presence of a stable isotope of chromium ( $\text{Cr}^{53}$ ) of about 9.55% abundance with electronic spin  $S = 3/2$  and nuclear spin  $I = 3/2$  [9]. The EPR spectrum is simulated using EasySpin [10] and the spin Hamiltonian parameters obtained as discussed later. The simulated EPR spectrum is shown in Fig. 1 (b).

EPR spectra of  $\text{Cr}^{3+}$  ion are analyzed using a spin Hamiltonian [9]:

$$\mathcal{H} = \mu_B B \cdot g \cdot S + D \left[ S_Z^2 - \frac{1}{3} S(S+1) \right] + E(S_X^2 - S_Y^2) + AIS \quad (1)$$

where the first, second, third and fourth terms represent the electron Zeeman interaction, axial zero-field splitting, rhombic zero-field splitting and hyperfine interaction, respectively. For a  $\text{Cr}^{3+}$  ( $d^3$ ) ion, the zero-field splitting arises from the interaction of the  $S = 3/2$  system with the non-cubic components of the ligand field.

The theoretical field positions of EPR transitions are obtained by using the equations of Ref. [11]. The calculated field positions of fine structure lines corresponding to different transitions for rotation in three mutually perpendicular planes are fitted to experimental values. The calculated (solid line) and experimental field positions (points) shown in Fig 2((a): site I; (b): site II) are in good agreement. The direction of the maximum overall splitting of EPR spectra is taken as the z-axis and that of the minimum as the x-axis [12]. The (x, y, z) system is parallel to the crystallographic **a**-, **b**-, **c**-axes. The Z-axis of symmetry adopted axes (SAA) is coincident with the crystal **c**-axis and the other two axes (X, Y) lie in the ab plane for site I, whereas the Z-axis of SAA is coincident with the crystal **b**-axis and the other two axes (X, Y) lie in the ac plane for site II. This indicates that the local site symmetry is approximately orthorhombic type I ( $D_{2h}$ ) [13]. The SAA system (X, Y, Z) of  $\text{Cr}^{3+}$  doped CTBZ single crystal for both sites is shown in Fig. 3. g-values and their direction cosines with respect to crystallographic axes are obtained using diagonalization procedure [14]. The best fit spin Hamiltonian parameters g, D, E and A are given in Table 1. The experimentally obtained direction cosines of g for  $\text{Cr}^{3+}$  ion (Table 1) are compared with the direction cosines of different bonds obtained from crystal structure data [8] shown in Table 2. The direction cosines of Cs-Br (1) and Cs-Br (3) bonds agree reasonably well with the experimental direction cosines of g of  $\text{Cr}^{3+}$  ions. This indicates that  $\text{Cr}^{3+}$  substitutes at  $\text{Cs}^+$  sites. Also the ionic radius of  $\text{Cr}^{3+}$  ( $0.63 \text{ \AA}$ ) is less than the ionic radius of  $\text{Cs}^+$  ( $1.74 \text{ \AA}$ ). Thus, the  $\text{Cr}^{3+}$  ion can fit well at the place of  $\text{Cs}^+$ . This supports the conclusion drawn on the basis of direction cosines. The charge compensation takes place by forming vacancies at some cation sites. The impurity ion enters the CTBZ lattice in a similar manner as observed in other  $\text{Cr}^{3+}$  doped crystal [15].

## 5. Optical Absorption Analysis

The electronic configuration of  $\text{Cr}^{3+}$  ion gives rise to the free ion terms  $^4F$ ,  $^4P$ ,  $^2G$  and several other doublet states out of which  $^4F$  is the ground state. In octahedral surroundings, the strong field configuration  $(t_{2g})^3$  yields the states  $^4A_{2g}$ ,  $^2E_g$ ,  $^2T_{1g}$  and  $^2T_{2g}$ . Out of these, the  $^4A_{2g}$  state lies lowest and is the ground state for  $\text{Cr}^{3+}$  ion. The first excited state configuration  $(t_{2g})^2(e_g)^1$  provides quartet states  $^4T_{1g}(F)$  and  $^4T_{2g}(F)$  and a number of doublet states, and in weak field  $^4F$  state splits as  $^4A_{2g}(F)$ ,  $^4T_{2g}(F)$  and  $^4T_{1g}(F)$  whereas  $^4P$  as  $^4T_{1g}(P)$ . In the doublet states,  $^2G$  splits as  $^2A_{1g}(G)$ ,  $^2E_g(G)$ ,  $^2T_{1g}(G)$  and  $^2T_{2g}(G)$  and  $^2H$  as  $^2E_g(H)$ ,  $^2T_{1g}(H)$  and  $^2T_{2g}(H)$ . The weak field terms  $^4A_{2g}(F)$ ,  $^2E_g(G)$ ,  $^2T_{1g}(G)$  and  $^2T_{2g}(G)$  correspond to the lowest strong field configuration  $t_{2g}^3$ . However, the ground state is  $^4A_{2g}(F)$  at all strengths of the crystal field. Thus, only two spin

allowed transitions are obtained from  ${}^4A_{2g}(F) \rightarrow {}^4T_{2g}(F)$  and  ${}^4A_{2g}(F) \rightarrow {}^4T_{1g}(F)$  states in addition to several spin forbidden transitions. If the symmetry is lowered, the degeneracy will be lifted and several absorption bands observed. The optical spectra of  $Cr^{3+}$  doped CTBZ recorded at room temperature are shown in Fig.4 (a, b). These exhibit two intense bands located at 18691 and 25000  $cm^{-1}$  and nine weak bands located at 11049, 12048, 13793, 15503, 17897, 20408, 31746, 40816 and 47619  $cm^{-1}$ , respectively. The two intense bands observed at 18,690 and 25,000  $cm^{-1}$  correspond to the transitions  ${}^4A_{2g}(F) \rightarrow {}^4T_{2g}(F)$  and  ${}^4A_{2g}(F) \rightarrow {}^4T_{1g}(F)$ , respectively. The three weak bands observed at 12,048, 15,503, 31,746  $cm^{-1}$  are assigned to the spin forbidden transitions from  ${}^4A_{2g}(F)$  to  ${}^2E_g(G)$ ,  ${}^2T_{1g}(G)$ ,  ${}^2A_{1g}(G)$ , respectively. The band at 11049  $cm^{-1}$  is the component band of the transition  ${}^4A_{2g}(F) \rightarrow {}^2E_g(G)$ . The two weak bands observed at 40,816 and 47619  $cm^{-1}$  are assigned to charge transfer (CT) bands [16]. Similarly, the band at 13793  $cm^{-1}$  is the component band of the transition  ${}^4A_{2g}(F) \rightarrow {}^2T_{1g}(G)$ . The bands at 17897 and 20408  $cm^{-1}$  are component bands of the transition  ${}^4A_{2g}(F) \rightarrow {}^4T_{2g}(F)$ .

The value of the crystal field parameter Dq and Racah parameter B are evaluated from the observed energy bands [17, 18]. The intense band position  $\nu_1$  at 18,691  $cm^{-1}$  corresponding to the transition  ${}^4A_{2g}(F) \rightarrow {}^4T_{2g}(F)$  gives the value 10Dq. The band at 25,000  $cm^{-1}$  corresponding to  ${}^4A_{2g}(F) \rightarrow {}^4T_{1g}(F)$  transition gives  $\nu_2$  and the Racah parameter B is given by the equation [19],

$$B = (2\nu_1^2 + \nu_2^2 - 3\nu_1\nu_2) / (15\nu_2 - 27\nu_1) \quad (2)$$

The value of B is evaluated as 602  $cm^{-1}$ , which is significantly smaller than the free ion value of  $Cr^{3+}$ ,  $B_{free} = 918 cm^{-1}$  [20, 21]. This decrease in B from  $B_{free}$  is caused by the effect of bond covalency. The Racah parameter C can be calculated from the position of the  ${}^4A_{2g}(F) \rightarrow {}^2E_g(G)$  transition using the following equation [22],

$$\frac{C}{B} = \frac{1}{3.05} \left[ \left( \frac{E(^2E)}{B} \right) - 7.9 + 1.8 \left( \frac{B}{Dq} \right) \right] \quad (3)$$

The energy matrices [17, 18] for  $d^3$  configuration in octahedral symmetry have been diagonalised for different values of B, C and Dq. The best fit of the observed bands is obtained for Dq = 1870, B = 602 and C = 2504  $cm^{-1}$ .

The calculated and experimental band positions are in good agreement for Dq = 1870  $cm^{-1}$  as shown in Fig 5. The band positions and their assignments are given in Table 3. The bonding in  $Cr^{3+}$ : CTBZ can be predicted by nephelauxetic parameters h and k of ligands and the central metal ion, respectively. h is obtained from the expression [23],

$$h = \left[ (B_{free} - B) / B_{free} \right] / k_{Cr^{3+}} \quad (4)$$

which is 1.63, since the value of k for  $Cr^{3+}$  is 0.21 [24]. The larger value of h gives delocalization of the d electrons suggesting the predominantly ionic nature of bonding between  $Cr^{3+}$  and ligands. By correlating optical and EPR data, the chemical bonding parameter  $\alpha'$  is obtained by [25],

$$g = g_e - 8\alpha' \lambda / \Delta \quad (5)$$

$$\lambda_{eff} = \alpha' \lambda$$

where  $g_e$  is the g-value of free electron ( 2.0023),  $\lambda$  is the spin-orbit coupling constant (91  $cm^{-1}$ ) [17],  $\Delta$  is the gap between the excited and ground levels,  $\lambda_{eff}$  is the spin-orbit coupling constant for ion ( $Cr^{3+}$ ) in the crystal [8]. From the experimental values  $g = 2.0104$  and  $\Delta = 18690 cm^{-1}$  for site I and  $g = 2.0061$  and  $\Delta = 18690 cm^{-1}$  for site II, the values of  $\alpha'$  are 0.21 for site I and 0.10 for site II.

The value  $0 < \alpha' < 1$  is characterized by the ionic contribution of the chemical bond between the  $Cr^{3+}$  ion and its ligands. Thus, the evaluated value of  $\alpha'$  suggests the ionic nature of  $Cr^{3+}$  ion in CTBZ. The value of  $\lambda_{eff}$  is evaluated as 19  $cm^{-1}$ (site I) and 9  $cm^{-1}$ (site II). Its deviation from the free ion value can be attributed to bonding effects in the crystal which yield a reduction in the effective value of spin-orbit coupling and also in orbital moment. The value of Dq/B obtained (3.1) is of the order of Dq/B = 2.8 found for  $Al_2O_3: Cr^{3+}$  [26], which shows that the  $Cr^{3+}$  ions in CTBZ are situated in a strong crystal field. The  ${}^2E_g$  state is unsplit in undistorted octahedral symmetry of the  $Cr^{3+}$  site [27]. But it splits in trigonal, tetragonal or orthorhombic distortions (low-symmetry fields) [21].

In the present study since the band associated with  ${}^2E_g$  shows splitting, the distortion can be of any of the above three types. The  $T_{1g}$  or  $T_{2g}$  level splits into two components each in tetragonal and three components each in orthorhombic distortion. Since the band associated with  $T_{2g}$  level here splits into three components, the distortion is expected to be of orthorhombic type [21]. The values of Dq, B and C are in the range reported for  $Cr(NH_3)_6^+$  [28]. We would expect the crystal field strength to be larger in CTBZ as in MgO [9] and therefore a Dq value as in  $Cr(NH_3)_6^+$  or even larger can be expected.

The spin-orbit splitting can be evaluated from the relation of Cole and Garret [29],

$$\lambda_{eff} = 0.0011(B + 1.08)3 + 0.0062 \quad (6)$$

where the Racah parameter B and spin-orbit coupling constant  $\lambda_{eff}$  are in kK. The value of  $\lambda_{eff}$ , thus determined from the above expression is  $79 \text{ cm}^{-1}$ . This value is in good agreement with the value expected for the  $Cr^{3+}$  ion in crystals having oxygen ligands, like in MgO [28]. It is interesting to compare this value with the value obtained from EPR and optical studies.

## 6. Theoretical Investigation

The energy level of the ground state of transition metal ion doped in crystal can be described using the spin Hamiltonian consisting of Zeeman electronic (Ze) and ZFS terms [30]. Explicit form of the ZFS terms for  $Cr^{3+}$  ion with  $S = 3/2$  at orthorhombic symmetry site ( $D_{2h}$ ) can be written as:

$$\begin{aligned} \mathcal{H} &= B_2^0 O_2^0 + B_2^2 O_2^2 + AIS = \frac{1}{3} b_2^0 O_2^0 + \frac{1}{3} b_2^2 O_2^2 + AIS \\ &= D(S_z^2 - \frac{1}{3} S(S+1)) + E(S_x^2 - S_y^2) + AIS \end{aligned} \quad (7)$$

The ZFSPs in Eq. (7) can be obtained using SPM. Conventional zero field splitting parameters were derived [31-34] for  $3d^3$  ion at orthorhombic symmetry as

$$D = b_2^0 = \frac{\bar{b}_2(R_0)}{2} \left[ \left( \frac{R_0}{R_i} \right)^{t_2} \sum_i (3 \cos^2 \theta_i - 1) \right] \quad (8)$$

$$E = \frac{b_2^2}{3} = \frac{\bar{b}_2(R_0)}{2} \left[ \left( \frac{R_0}{R_i} \right)^{t_2} \sum_i \sin^2 \theta_i \cos 2\phi_i \right] \quad (9)$$

For  $Cr^{3+}$  ions, the parameters of the rank  $k = 2$  and 4 exist for crystal field parameters (CFPs), while only of the rank  $k = 2$  for ZFSPs.

In CTBZ single crystal, the local symmetry around  $Cr^{3+}$  ion is slightly lower than orthorhombic. However, the calculations are done under the approximation of orthorhombic symmetry. For octahedral co-ordination around  $Cr^{3+}$  ion in the crystal, the parameters  $t_2$  and  $\bar{b}_2(R_0)$  for  $Cr^{3+}$  ion in case of  $LiNbO_3$  having  $Cr^{3+}-O^{2-}$  bond are:  $t_2 = 0.12$  and  $\bar{b}_2(R_0) = 2.34 \text{ cm}^{-1}$  [35]. The above parameters are used in our case for  $Cr^{3+}-O^{2-}$  bond.

$Cr^{3+}$  ion enters the lattice of CTBZ substitutionally at site I and site II. The position of metal ion and spherical coordinates of ligands are given in Table 4. The calculated and conventional ZFS parameters together with reference distance are shown in Table 5. For the evaluation of ZFSPs, the reference distances of 0.205 nm and 0.179 nm were fixed. The calculated conventional ZFSPs are:  $D = 240 \times 10^{-4} \text{ cm}^{-1}$ ,  $E = 64 \times 10^{-4} \text{ cm}^{-1}$  and  $D = 230 \times 10^{-4} \text{ cm}^{-1}$ ,  $E = 75 \times 10^{-4} \text{ cm}^{-1}$  for site I and site II, respectively. The ratio  $b_2^2/b_2^0$  is within the range (0, 1) for orthorhombic symmetry [12]. Here, the ratio of  $b_2^2/b_2^0 = 0.800$ ,  $E/D = 0.266$  for site I and  $b_2^2/b_2^0 = 0.937$ ,  $E/D = 0.326$  for site II. The theoretical ZFS parameters are in reasonable agreement with the experimental ones. The other reference distance obtained for oxide ligands [30] was also used but the ZFSPs determined were found very different as compared to the experimental values.

The CFPs for  $Cr^{3+}$  in crystals are obtained by the SPM formula [36]:

$$B_{kq} = \sum_i \bar{B}_k \left( \frac{R_0}{R_i} \right)^{t_k} K_{kq}(\theta_i, \phi_i) \quad (11)$$

where  $R_0$  is the reference distance for the site;  $R_i$ ,  $\theta_i$ ,  $\phi_i$  are the polar coordinates of the  $i^{\text{th}}$  ligand and  $K_{kq}$  is the coordination factor [37]. For calculating  $B_{kq}$  ( $k = 2, 4$ ;  $q = 0, 2, 4$ ),  $\bar{B}_2 = 40, 400 \text{ cm}^{-1}$ ,  $t_2 = 1.3$ ,  $\bar{B}_4 = 11, 700 \text{ cm}^{-1}$  and  $t_4 = 3.4$  are taken from Ref. [37]. The calculated  $B_{kq}$  parameters are given in Table 6. These parameters are used to calculate optical absorption spectra with the help of CFA program [38-39]. The energy levels of the impurity ion are obtained by diagonalizing the complete Hamiltonian within the  $3d^N$  basis of states in the intermediate crystal field coupling scheme. The Hamiltonian contains the Coulomb interaction (in terms of B and C), Trees correction, the spin-orbit interaction, the crystal field Hamiltonian, the spin-spin interaction and the spin-other-orbit interaction. The calculated energy values are given in Table 7 together with the experimental values for comparison. It is found that there is good agreement between the two values.

## 7. Conclusion

EPR and optical studies of  $\text{Cr}^{3+}$  doped CTBZ single crystal have been carried out at room temperature. The spin Hamiltonian parameters (D, E, g and A) have been evaluated. The  $\text{Cr}^{3+}$  ions enter the lattice substitutionally by replacing  $\text{Cs}^+$  ions. On the basis of EPR and optical data, the nature of bonding in the crystal has also been discussed. The values of the bonding parameters both from optical and EPR studies suggest that the bonding between  $\text{Cr}^{3+}$  and ligands is ionic. The hyperfine structure has been observed due to interaction of unpaired electron spin and nuclear spin.

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### Table and Figure Captions:

Table 1. Principal g values, their direction cosines and the parameters D, E and A.

Table 2. The ionic distances and direction cosines of different bonds of Cr<sup>3+</sup> doped CTBZ single crystal.

Table 3. Observed and calculated energy bands and their assignments for Cr<sup>3+</sup> doped CTBZ (Dq = 1870 cm<sup>-1</sup>, B = 602 cm<sup>-1</sup> and C = 2504 cm<sup>-1</sup>).

Table 4. Fractional positions of Cr<sup>3+</sup> ions together with spherical co-ordinates (R,  $\theta$ ,  $\phi$ ) of ligands in CTBZ single crystals.

Table 5. Calculated and conventional zero field splitting parameters together with Reference distance for Cr<sup>3+</sup> doped CTBZ single crystal.

Table 6. B<sub>kq</sub> parameters to be used in CFA program for calculating optical absorption spectra.

Table 7. Observed and calculated energy bands and their assignments for Cr<sup>3+</sup> doped CTBZ single crystal using CFA program.

Fig. 1. (a) EPR spectrum recorded in ab plane for magnetic field B at 0<sup>0</sup> from c axis.

(b) Simulated EPR spectrum in ab plane for magnetic field B at 0<sup>0</sup> from c axis (Frequency 9.11 GHz).

Fig. 2. Angular variation of EPR lines for Cr<sup>3+</sup> doped CTBZ in all three planes for (a) Site I, (b) Site II.

Fig. 3. Crystal structure of CTBZ in which symmetry adopted axis system (SAAS) is shown.



Fig. 4 . Optical absorption spectrum of Cr<sup>3+</sup> doped CTBZ recorded at room temperature in the wavelength range (a) 195-25 nm, (b) 325-1100 nm.

Fig. 5. Variation of energy levels of Cr<sup>3+</sup> doped CTBZ with Dq for B = 602 cm<sup>-1</sup>, C = 2504 cm<sup>-1</sup> and α = 76 cm<sup>-1</sup> (experimental energies are shown by circles).

**Table 1**

Site	Principal g values	a	b	c	D	E	A
I	g <sub>x</sub> = 1.9831	-0.9900	0.0109	0.1403	234	69	80
	g <sub>y</sub> = 2.0147	0.0242	0.9954	0.0930			
	g <sub>z</sub> = 2.0336	<b>0.1387</b>	<b>-0.0955</b>	<b>0.9857</b>			
II	g <sub>x</sub> = 1.9840	-0.9942	0.1007	-0.0379	235	70	82
	g <sub>y</sub> = 2.0128	0.1028	0.9931	-0.0563			
	g <sub>z</sub> = 2.0215	<b>-0.0320</b>	<b>0.0599</b>	<b>0.9977</b>			

D, E and A are in units of 10<sup>-4</sup> cm<sup>-1</sup> and estimated errors for g, D, E and A are ±0.0002 and ±2×10<sup>-4</sup> cm<sup>-1</sup> respectively.

**Table 2**

Bonds	Ionic Distances(Å)	Direction cosines with respect to crystallographic axes		
	R(Å)	a	b	c
Cs-Br <sup>-</sup> (1)	6.4728	±0.6422	±0	±0.7664
Cs-Br <sup>-</sup> (2)	4.2392	±0.4649	±0	±0.8853
Cs-Br <sup>-</sup> (3)	3.7918	±0.2208	±0.2785	±0.9346
Cs-Br <sup>-</sup> (4)	3.7918	±0.2208	±0.2785	±0.9346
Cs-Br <sup>-</sup> (5)	3.7918	±0.8397	±0.5122	±0.1800
Cs-Br <sup>-</sup> (6)	4.2650	±0.3222	±0	±0.9466

**Table 3**

Transition from	Observed wave number (cm <sup>-1</sup> )	Calculated wave number (cm <sup>-1</sup> )
<sup>4</sup> A <sub>2g</sub> (F)		
<sup>2</sup> E <sub>g</sub> (G)	11049, <b>12048</b>	12200
<sup>2</sup> T <sub>1g</sub> (G)	13793, <b>15503</b>	14992
<sup>4</sup> T <sub>2g</sub> (F)	17897, <b>18691</b> , 20408	18700
<sup>4</sup> T <sub>1g</sub> (F)	25000	25642
<sup>2</sup> A <sub>1g</sub> (G)	31746	30490
<sup>2</sup> T <sub>2g</sub> (H)	40816	39648
<sup>2</sup> T <sub>1g</sub> (H)	47619	47826

**Table 4**

Position of Fe <sup>3+</sup> (Fractional)	Ligands	Spherical co-ordinates of ligand		
		R(Å)	θ	φ
Site I : Substitutional ( 0,0,0)	Br <sup>-</sup> (1)	6.4728	90	0
	Br <sup>-</sup> (2)	4.2392	128	0
	Br <sup>-</sup> (3)	3.7918	79	31
	Br <sup>-</sup> (4)	3.7918	79	31
	Br <sup>-</sup> (5)	3.7918	79	-31
	Br <sup>-</sup> (6)	4.2650	18	0
Site II : Substitutional ( 0,0,0)	Br <sup>-</sup> (1)	6.8382	92	0
	Br <sup>-</sup> (2)	4.6435	94	0
	Br <sup>-</sup> (3)	3.1746	83	33
	Br <sup>-</sup> (4)	3.1746	83	33
	Br <sup>-</sup> (5)	3.2820	81	30
	Br <sup>-</sup> (6)	4.2910	24	0

**Table 5**

	R <sub>0</sub> (Å)	Zero-field splitting parameters (cm <sup>-1</sup> )			Conventional Zero-field splitting parameters (×10 <sup>-4</sup> cm <sup>-1</sup> )		
		$b_2^0$	$b_2^2$	$b_2^2 / b_2^0$	D	E	E/D
Site I	2.05	-0.0240	0.0192	0.800	240	64	0.266
Site II	1.79	-0.0230	0.0225	0.937	230	75	0.326



**Table 6**

<u><math>B_{kq}</math> (<math>\text{cm}^{-1}</math>) Parameters used for CFA package</u>						
	$R_0(\text{\AA})$	$B_{20}$	$B_{22}$	$B_{40}$	$B_{42}$	$B_{44}$
Site I	2.05	54103	18917	-24047	11723	-38461
Site II	1.79	56107	17915	-23049	10726	-39484

**Table 7**

Transition from $^4A_{2g}(F)$	Experimentally observed band ( $\text{cm}^{-1}$ )	Calculated energy bands from CFA ( $\text{cm}^{-1}$ )
$^2E_g(G)$	11049, <b>12048</b>	12342,12652
$^2T_{1g}(G)$	13793, <b>15503</b>	13731,14837,15362
$^4T_{2g}(F)$	17857, <b>18691</b> ,20408	17586,17439,17785,18261,18285,18430
$^4T_{1g}(F)$	25000	24088,24120,24180,24315,24344,24420
$^2A_{1g}(G)$	31746	31583
CT band	40816	40615
CT band	47619	47639

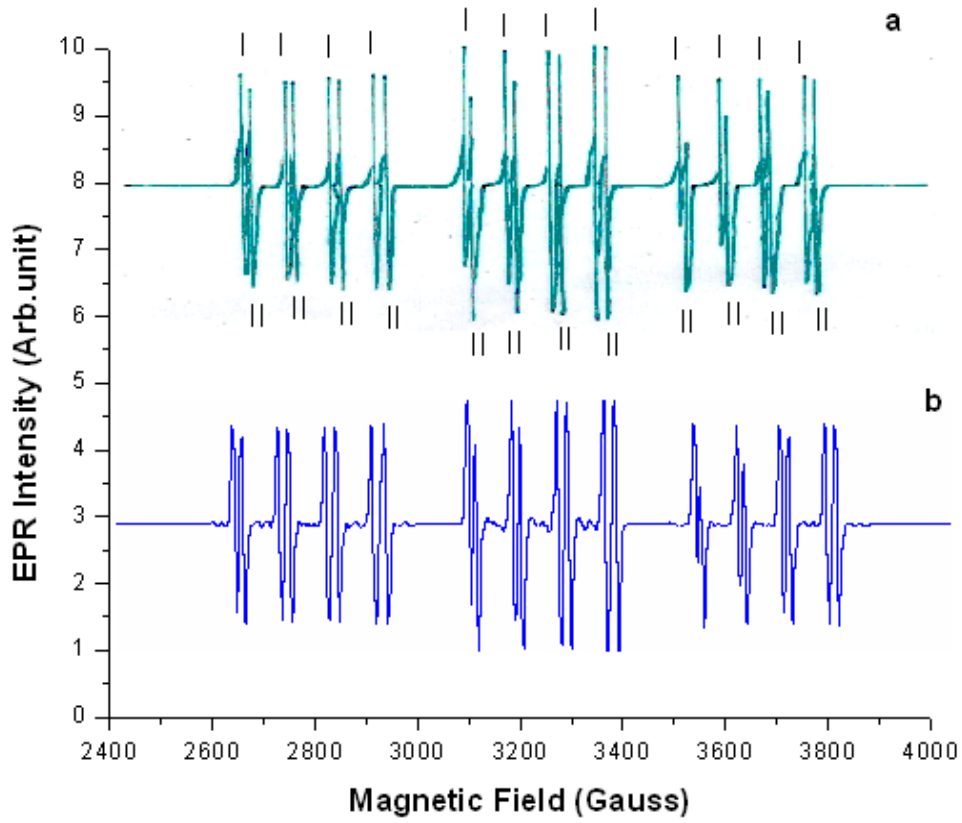


Fig. 1

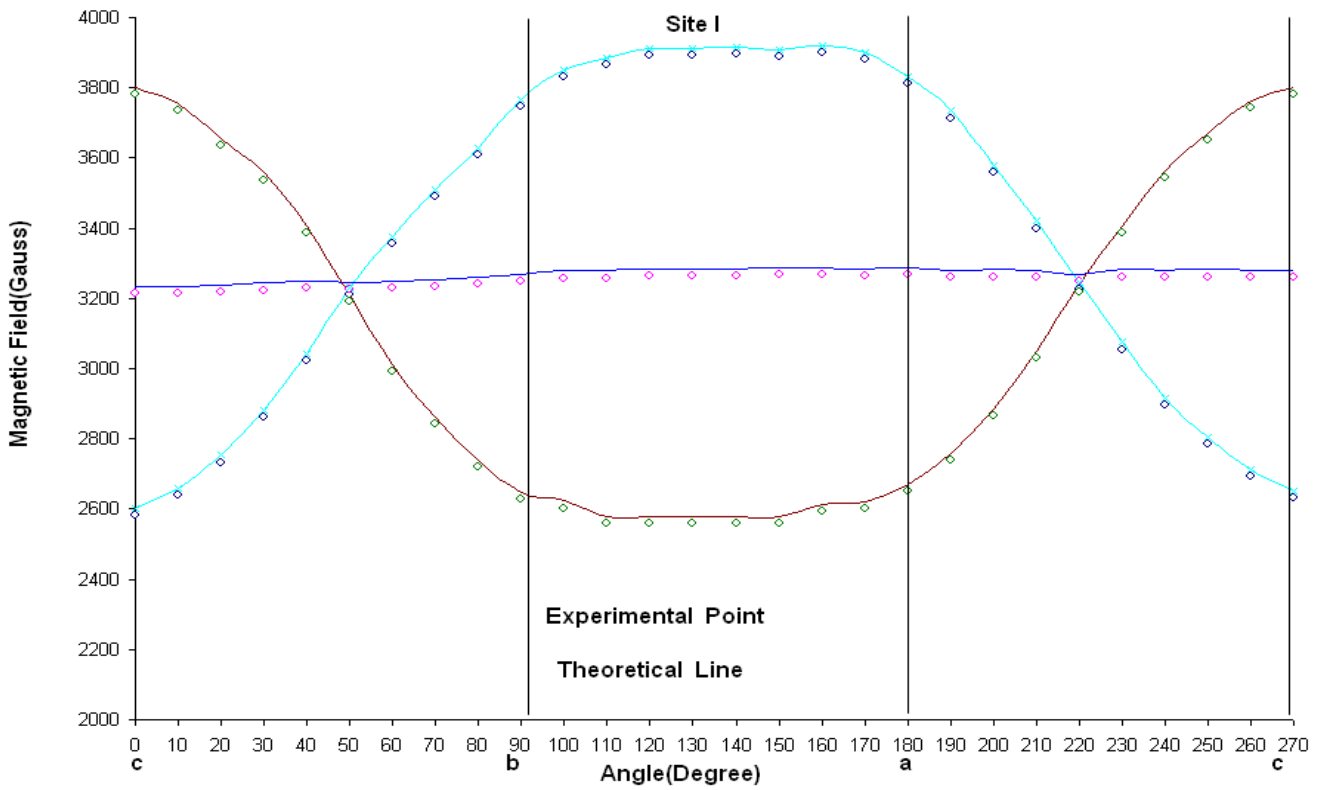


Fig. 2 (a)

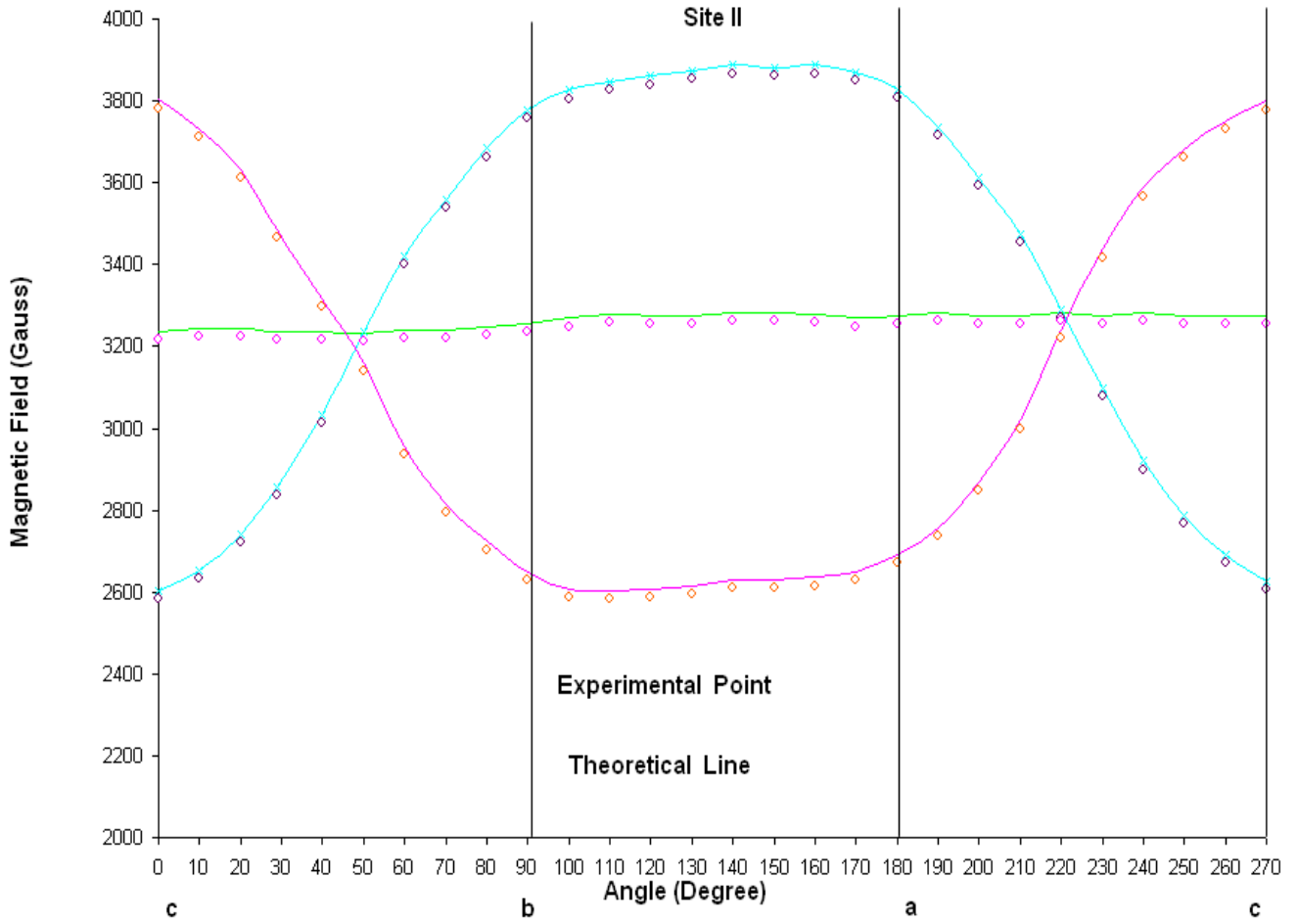


Fig. 2(b)

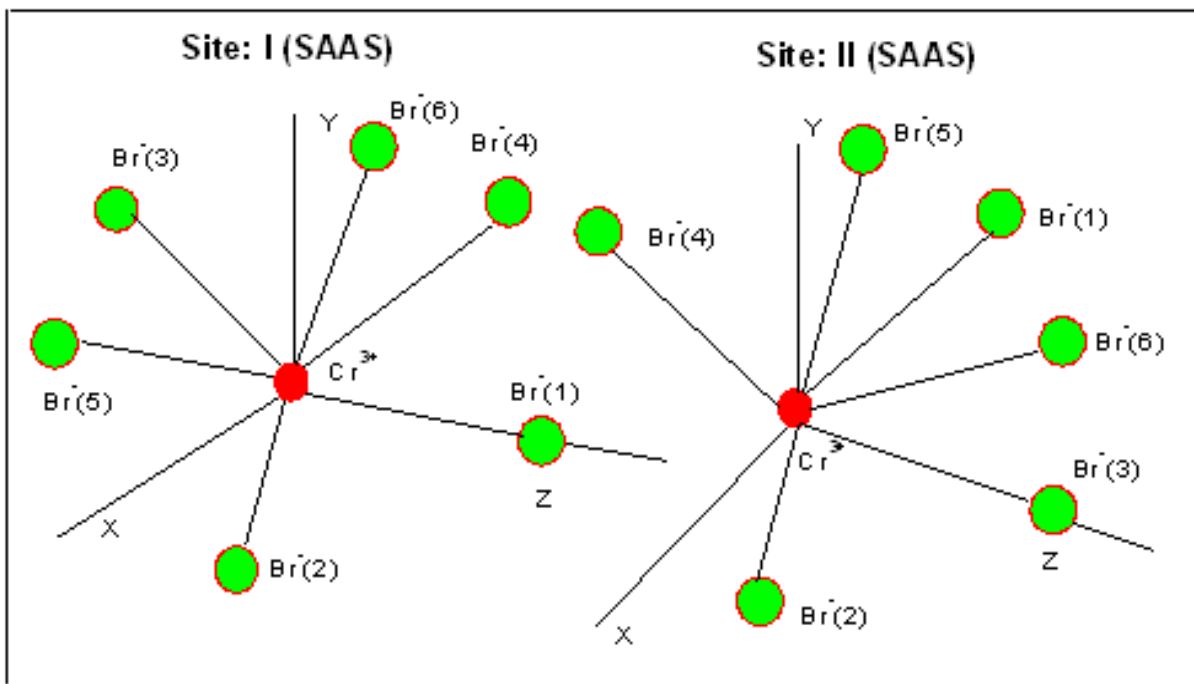


Fig. 3

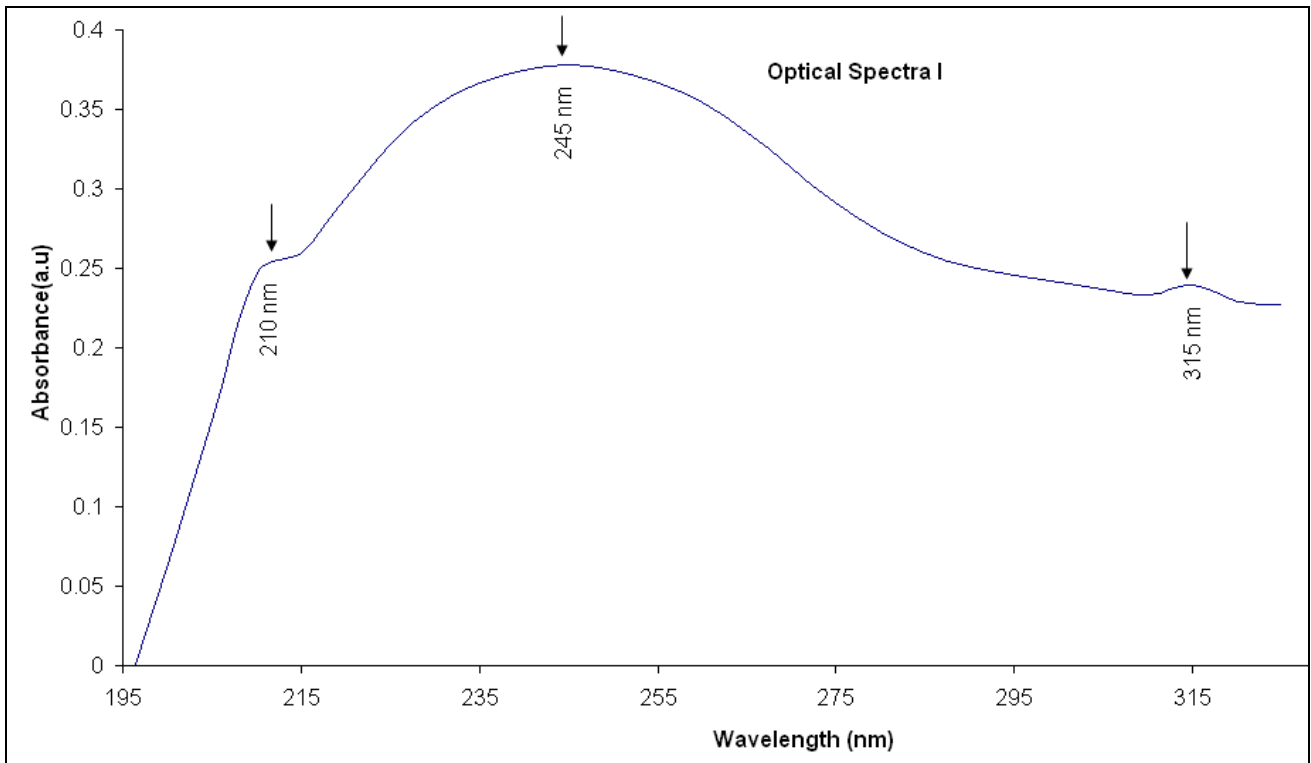


Fig. 4(a)

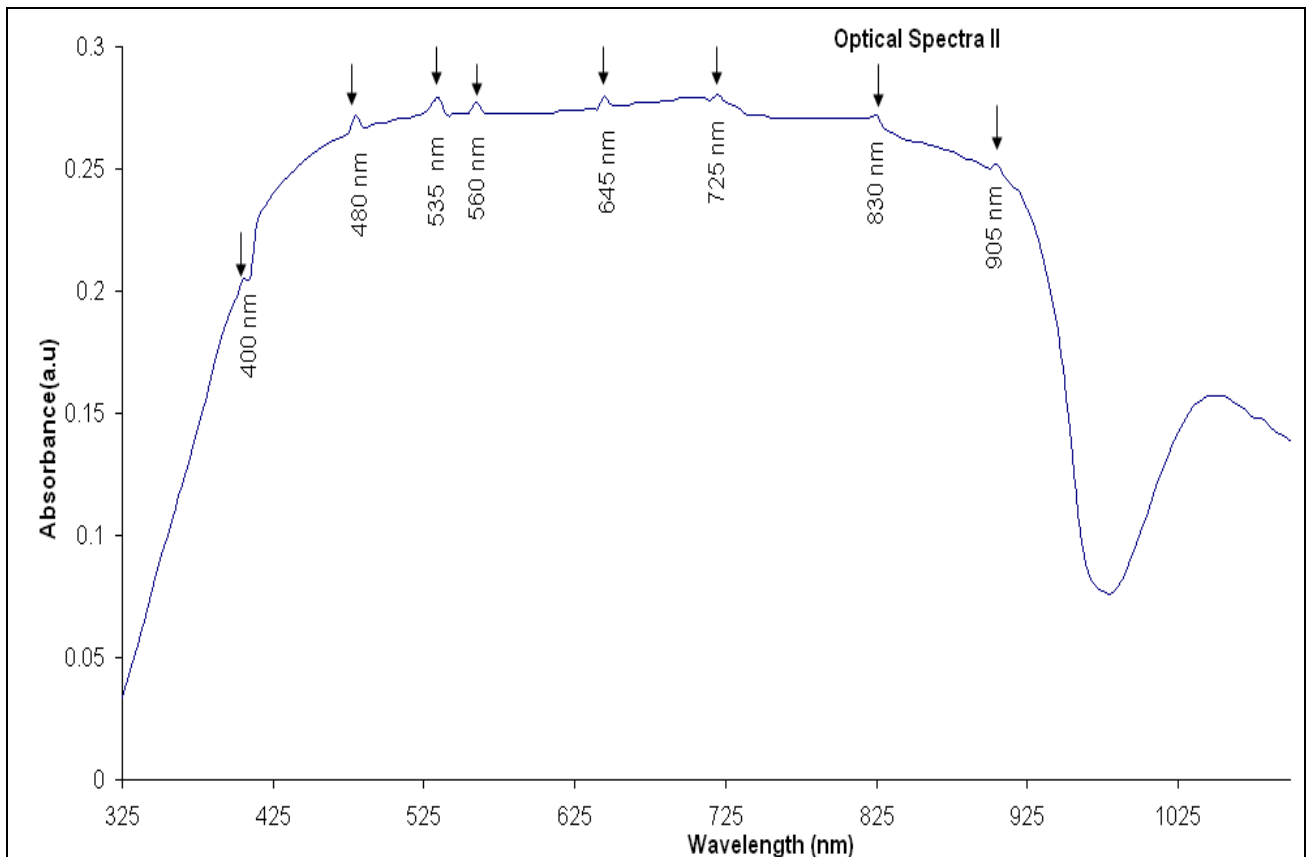


Fig. 4(b)

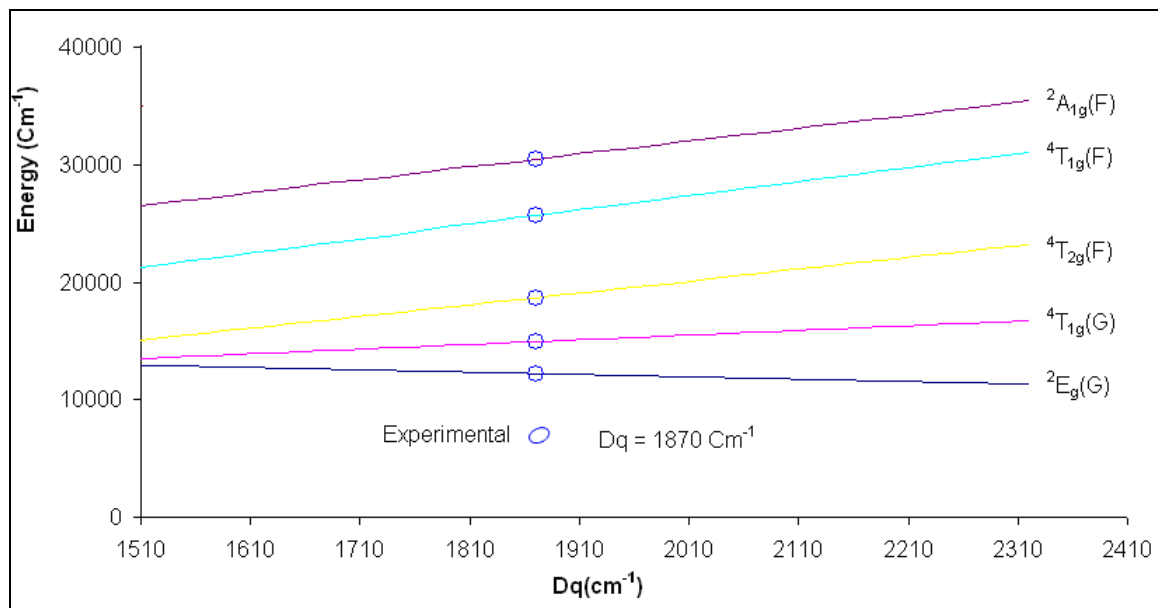


Fig. 5